

Gas Accretion onto a Supermassive Black Hole: a step to modeling AGN feedback in cosmological simulations

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# Outline

- Intro / Motivation
- The simplest case: spherical Bondi accretion
- Include radiative cooling / heating -- radiative feedback
   by X-rays
   Barai, Proga, KN, 2011, MNRAS, in press (arXiv:1102.3925)
- Non-spherical accretion flow, fragmentation due to thermal instability
   Barai, Proga, KN, 2011, in prep. (Paper II)
- Conclusions

## **Motivation**



- Still a large gap btw small-scale sims & cosmological sims. (≈ pc) (~kpc - 10 Mpc)
- Cosmo sims uses ad-hoc AGN accretion models as "sub-grid" physics.
- How well can a cosmological SPH code (e.g. GADGET) handle accretion onto a SMBH?

### **The Bondi Accretion Problem**

- Spherically symmetric accretion onto a central mass (Bondi 1952)
- Gas is at rest at infinity, with  $\rho_{\infty} \& p_{\infty}$ . Increase in the central mass is ignored.
- Two equations are solved:

$$\dot{M} = -4\pi r^2 \rho v = \text{constant.} \quad \text{(Continuity Eq.)}$$
$$\frac{v^2}{2} + \left(\frac{\gamma}{\gamma - 1}\right) \frac{p_{\infty}}{\rho_{\infty}} \left[ \left(\frac{\rho}{\rho_{\infty}}\right)^{\gamma - 1} - 1 \right] = \frac{GM_{BH}}{r}, \quad \text{(Bernoulli's Eq.)}$$

• One of the solutions:

$$\dot{M}_B = 4\pi\lambda_c \frac{(GM_{BH})^2}{c_{s,\infty}^3} \rho_{\infty}, \qquad \lambda_c = \left(\frac{1}{2}\right)^{\frac{(\gamma+1)}{2(\gamma-1)}} \left(\frac{5-3\gamma}{4}\right)^{\frac{(3\gamma-5)}{2(\gamma-1)}}.$$

• Characteristic scales:

**Bondi radius:** 
$$R_B = \frac{GM_{BH}}{c_{s,\infty}^2}$$
. **Sonic radius:**  $R_s = \left(\frac{5-3\gamma}{4}\right)R_B$ .  
**Bondi time:**  $t_B = \frac{R_B}{c_s} = \frac{GM_{BH}}{c_{s,\infty}^3}$ .

## <u>Simplest Case: Spherical Bondi</u> <u>Accretion Flow onto a SMBH</u>

- GADGET-3: 3-d cosmological SPH/ N-body code (Springel '05)
- Central SMBH 10<sup>8</sup> M<sub>☉</sub> represented by a pseudo-Newtonian Paczynsky & Wiita (1980) potential
- r<sub>out</sub>=5-20 pc, N<sub>ptcl</sub>=64<sup>3</sup>-128<sup>3</sup>
- Set IC to uniform/spherical Bondi flow w/  $\gamma$ =1.01,  $\rho_{\infty}$ =10<sup>-19</sup> g/cm<sup>3</sup>, T<sub>\infty</sub>=10<sup>7</sup>K, T<sub>init</sub>=T<sub>\infty</sub>
- Corresponding Bondi solution: R<sub>B</sub>=3pc, R<sub>sonic</sub>=1.5pc, t<sub>B</sub>=7.9e3yr
- All runs: **r**<sub>in</sub>=**0.1pc**, γ=1.01

## 3-d spherical volume, vacuum boundary condition

Run No.	$r_{ m out} \ [ m pc]$	N <sup>b</sup>	IC	${M_{ m tot,IC}}$ c $[M_{\odot}]$	${M_{ m part}}^{ m d}$ $[M_{\odot}]$	$t_{ m end}  {}^{ m e}$ $[10^4 { m yr}]$	
1	5	$64^{3}$	Uniform <sup>i</sup>	$3.96 \times 10^{5}$	1.51	3	
2	10	$64^{3}$	Uniform	$6.19 \times 10^6$	23.61	7.2	
3	50	$128^{3}$	Uniform	$7.73 \times 10^{8}$	368.60	20	
4	5	$64^{3}$	Bondi <sup>j</sup>	$1.81 \times 10^6$	6.89	2	
5	10	$64^{3}$	Bondi	$9.76 \times 10^{6}$	37.23	8	
6	10	$128^{3}$	Bondi	$9.76 \times 10^{6}$	4.65	8	
7	20	128 <sup>3</sup>	Bondi	$6.24 \times 10^{7}$	29.75	8	
1	20	120	Donar	0.24 × 10	20.10	0	
$7a^{\mathrm{k}}$	20	$128^{3}$	Bondi	$6.24 \times 10^7$	29.75	80	
$7b^{-1}$	20	$128^{3}$	Bondi	$6.24 \times 10^7$	29.75	100	
8	50	$128^{3}$	Bondi	$8.48 \times 10^{8}$	404.35	16	
9	20	$128^{3}$	$\rho_B, v_{\text{init}} = 0$	$6.24 \times 10^{7}$	29.75	8	
10	20	$128^{3}$	Uniform	$4.95 \times 10^{7}$	23.60	8	
11	20	$128^{3}$	Hernquist $^{\rm m}$	$6.24 \times 10^7$	29.75	7.2	
12 <sup>n</sup>	20	$128^{3}$	Bondi	$6.24 \times 10^{7}$	29.75	8	

## **Example: Properties of Particles**



## Mass Inflow Rates at rin

- the larger r<sub>out</sub>, the longer duration of Bondi inflow rate
- If started from a Bondi flow, Bondi rate is achieved quickly.
- After a while, the inflow rate decreases due to the artificial outflow at the outer boundary.
- Greater sim. volume reduces this effect on mass inflow.

cm°3

Density (g





## **Radiative Heating & Cooling**

• Xray emitting corona irradiates the accretion flow

$$L_X = f_X L_{\text{Edd}}, \qquad L_{\text{Edd}} = \frac{4\pi c G m_p M_{BH}}{\sigma_e}, \qquad \qquad \mathbf{Flux:} \quad F_X = \frac{L_X}{4\pi r^2}.$$

 Approx. analytic heating/cooling rates from Blondin '94; opt-thin gas illuminated by a 10 keV bremsstrahlung.

net rate: $\rho \mathcal{L} = n^2 (G_{\text{Compton}} + G_X - L_{b,l})$  [erg cm<sup>-3</sup> s<sup>-1</sup>],Compton h/c rate: $G_{\text{Compton}} = 8.9 \times 10^{-36} \xi (T_X - 4T)$  [erg cm<sup>3</sup> s<sup>-1</sup>].Net Xray photoioniz. heating<br/>and recomb. cooling rate: $G_X = 1.5 \times 10^{-21} \xi^{1/4} T^{-1/2} \left(1 - \frac{T}{T_X}\right)$  [erg cm<sup>3</sup> s<sup>-1</sup>].Brems. and line cooling rate: $L_{b,l} = 3.3 \times 10^{-27} T^{1/2}$ <br/> $+ [1.7 \times 10^{-18} \exp(-1.3 \times 10^5/T) \xi^{-1} T^{-1/2}$ <br/> $+ 10^{-24}] \delta$  [erg cm<sup>3</sup> s<sup>-1</sup>].

 $T_X=1.16 \times 10^8 \text{ K}$  (=10keV, Blondin '94)

## Runs with radiative cooling/heating

Run No.	$r_{ m out} \ [ m pc]$	N	$M_{ m tot,IC}$ $[M_{\odot}]$	$M_{ m part}$ $[M_{\odot}]$	$\gamma_{ m init}$	$T_{\infty}$ [K]	$R_B$ [pc]	$ ho_{\infty}$ [g/cm <sup>3</sup> ]	$T_{\mathrm{init}}$	$L_X$ $[L_{ m Edd}]$	$t_{ m end} \ [10^5 \ { m yr}]$
$\begin{array}{c} 13 \\ 14 \end{array}$	$\begin{array}{c} 20\\ 50 \end{array}$	$\frac{128^3}{128^3}$	$5.81 \times 10^{5}$ $8.23 \times 10^{6}$	$\begin{array}{c} 0.277\\ 3.92 \end{array}$	$\begin{array}{c} 1.4 \\ 1.4 \end{array}$	$\frac{10^7}{10^7}$	$\begin{array}{c} 2.19 \\ 2.19 \end{array}$	$10^{-21}$ $10^{-21}$	$T_{\infty}$ $T_{\infty}$	$\begin{array}{c} 0.5 \\ 0.5 \end{array}$	$1.0\\2.9$
$15\\16$	20 20	$\frac{128^3}{256^3}$	$5.81 \times 10^{-1}$ $5.81 \times 10^{-1}$	$2.77 \times 10^{-7}$ $3.46 \times 10^{-8}$	$\begin{array}{c} 1.4 \\ 1.4 \end{array}$	$\frac{10^7}{10^7}$	$2.19 \\ 2.19$	$10^{-27}$ $10^{-27}$	$T_{\infty}$ $T_{\infty}$	$\begin{array}{c} 0.5\\ 5\times 10^{-4}\end{array}$	$\begin{array}{c} 1.0\\ 1.9\end{array}$
17 $18$ $19$	20 20 20	$128^3$ $128^3$ $128^3$	$5.81 \times 10^{5}$ $5.65 \times 10^{5}$ $1.47 \times 10^{7}$	$0.277 \\ 0.269 \\ 7.0$	$1.4 \\ 5/3 \\ 5/3$	$10^{7}$ $10^{7}$ $10^{5}$	2.19 1.84 183.9	$10^{-21}$ $10^{-21}$ $10^{-21}$	${T_{ m rad}}^{ m b}$ ${T_{ m rad}}$ ${T_{ m rad}}$	$5 \times 10^{-4}$ $5 \times 10^{-4}$ $5 \times 10^{-4}$	$2.9 \\ 3.0 \\ 1.5$
20 21	$200\\200$	$256^{3}$ $256^{3}$	$1.33 \times 10^9$ $4.95 \times 10^8$	79.09 29.50	5/3 5/3	$10^{5}$ $10^{7}$	$183.9\\1.84$	$10^{-21}$ $10^{-21}$	$T_{ m rad}$ $T_{ m rad}$	$5 \times 10^{-4}$ $5 \times 10^{-4}$	$\begin{array}{c} 6.5 \\ 8.7 \end{array}$
22 23	200 200	$\frac{128^3}{256^3}$	$1.33 \times 10^{7}$ $1.33 \times 10^{7}$	6.33 0.791	$\frac{5/3}{5/3}$	$\frac{10^{5}}{10^{5}}$	$\frac{183.9}{183.9}$	$10^{-23}$ $10^{-23}$	$\frac{T_{\rm rad}}{T_{\rm rad}}$	$\frac{5 \times 10^{-4}}{5 \times 10^{-4}}$	70 20
24 <sup>c</sup> 25 26	$200 \\ 200 \\ 200 \\ 200$	$1.24 \times 10^{7}$ $1.24 \times 10^{7}$ $1.24 \times 10^{7}$	$9.77 \times 10^{6}$ $9.77 \times 10^{6}$ $9.77 \times 10^{6}$	0.791 0.791 0.791	5/3 5/3 5/3	$10^{5}$ $10^{5}$ $10^{5}$	183.9 183.9 183.9	$10^{-23}$ $10^{-23}$ $10^{-23}$	$T_{Run23}$ $T_{Run23}$ $T_{Run22}$	$5 \times 10^{-5}$ $5 \times 10^{-3}$ $1 \times 10^{-2}$	19 21 22
20 27 28	200 200	$1.24 \times 10^{7}$ $1.24 \times 10^{7}$ $1.24 \times 10^{7}$	$9.77 \times 10^{6}$ $9.77 \times 10^{6}$	0.791 0.791	5/3 5/3	$\frac{10^5}{10^5}$	183.9 183.9	$10^{-23}$ $10^{-23}$	$T_{ m Run23}$ $T_{ m Run23}$ $T_{ m Run23}$	$ \begin{array}{c} 1 \times 10 \\ 2 \times 10^{-2} \\ 5 \times 10^{-2} \end{array} $	$25 \\ 50$

#### **Ptcl properties w/ radiative heating & cooling**

#### Representative run:

- red: free-fall scaling
- blue: ZEUS-2d result
- Near the inner radius, excess heating by artificial viscosity is seen.
- Inflow rate is enhanced above Bondi rate, due to lower gas temp: T(r<sub>out</sub>) <10<sup>5</sup> K, T∞=10<sup>5</sup> K





green: free-fall scaling w/ only adiabatic term

T<sub>ff,ar</sub>: solving internal energy eq. w/ both radiative & adiabatic term

## Impact of varying Lx on inflow rates

- Restart Run #23 at t=1.4 Myr, Lx/L<sub>Edd</sub>=5e-4 orig.
- Runs 24-28: increase Lx
- Dramatic decrease in M<sub>in</sub> at Lx/L<sub>Edd</sub>>0.01 ---transition from net inflow to net outflow



net outflow; non-spherical fragmentation observed.

#### Thermal instability due to rad. feedback

# **Non-spherical outflow:** due to rad. feedback Run 26: r<sub>out</sub>=200pc, Lx/L<sub>Edd</sub>=0.01



### **Ptcl Properties: impact of rad feedback**



## **Time Evolution of a Single Ptcl**

Run 26: r<sub>out</sub>=200pc, Lx/L<sub>Edd</sub>=0.01, t=2.0 Myr

- Start (triangle): r=53 pc, t=1.4 Myr
- End (square): r=lpc, t=1.8 Myr
- + symbol: dt=0.004 Myr



# **Non-spherical outflow:** Run 27: r<sub>out</sub>=200pc, Lx/L<sub>Edd</sub>=0.02 due to rad. feedback



### **Ptcl Properties: impact of rad feedback**



Run 27: r<sub>out</sub>=200pc, Lx/L<sub>Edd</sub>=0.02

# **Non-spherical outflow:** Run 28: r<sub>out</sub>=200pc, Lx/L<sub>Edd</sub>=0.05 due to rad. feedback

#### Temperature

#### Density



#### ± 200pc

## **Run 28**

#### r<sub>out</sub>=200pc, Lx/L<sub>Edd</sub>=0.05



log gas temperature

# <u>Conclusions</u>

- GADGET-3 SPH code can reproduce the spherical Bondi accretion rate properly, but with some limitations.
- spurious heating by Artificial Viscosity near r<sub>in</sub> & artificial outflow at r<sub>out</sub> due to outer BC are problems for SPH.
- non-spherical in/outflow develops due to rad. feedback via thermal instability, even in the simplest situation that we studied --- connection with NLR? (Paper II)
- Future work: include rad. pressure, rotation, diff geometry, comparison w/ NLR obs, connect with cosmological sim





± 30 pc range (t = 2.047 Myr)

colder, denser filament-like structures due to non-spherical fragmentation

#### Run 26: r<sub>out</sub>=200pc, Lx/L<sub>Edd</sub>=0.01

#### Zoom-in: inner 4 pc



