

Gas Accretion onto a Supermassive Black Hole: a step to modeling AGN feedback in cosmological simulations

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\begin{array}{r}
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\end{array}
$$

## Outline

- Intro / Motivation
- The simplest case: spherical Bondi accretion
- Include radiative cooling / heating -- radiative feedback by X-rays

Barai, Proga, KN, 20 I I, MNRAS, in press (arXiv: I I 02.3925)

- Non-spherical accretion flow, fragmentation due to thermal instability

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\text { Barai, Proga, KN, } 201 \text { I, in prep. (Paper II) }
$$

- Conclusions


## Motivation

Small-scale sims


Cosmological sims

(e.g. Di Matteo+, Booth \& Schaye, ...)

- Still a large gap btw small-scale sims \& cosmological sims.

$$
(\lesssim \mathrm{pc})
$$

(~kpc - 10 Mpc )

- Cosmo sims uses ad-hoc AGN accretion models as "sub-grid" physics.
- How well can a cosmological SPH code (e.g. GADGET) handle accretion onto a SMBH?


## The Bondi Accretion Problem

- Spherically symmetric accretion onto a central mass (Bondi 1952)
- Gas is at rest at infinity, with $\rho_{\infty} \& p_{\infty}$. Increase in the central mass is ignored.
- Two equations are solved:

$$
\begin{gathered}
\dot{M}=-4 \pi r^{2} \rho v=\text { constant. (Continuity Eq.) } \\
\frac{v^{2}}{2}+\left(\frac{\gamma}{\gamma-1}\right) \frac{p_{\infty}}{\rho_{\infty}}\left[\left(\frac{\rho}{\rho_{\infty}}\right)^{\gamma-1}-1\right]=\frac{G M_{B H}}{r}, \quad \text { (Bernoulli's Eq.) }
\end{gathered}
$$

- One of the solutions:

$$
\dot{M}_{B}=4 \pi \lambda_{c} \frac{\left(G M_{B H}\right)^{2}}{c_{s, \infty}^{3}} \rho_{\infty}, \quad \quad \lambda_{c}=\left(\frac{1}{2}\right)^{\frac{(\gamma+1)}{2(\gamma-1)}}\left(\frac{5-3 \gamma}{4}\right)^{\frac{(3 \gamma-5)}{2(\gamma-1)}} .
$$

- Characteristic scales:

Bondi radius: $\quad R_{B}=\frac{G M_{B H}}{c_{s, \infty}^{2}} . \quad$ Sonic radius: $R_{s}=\left(\frac{5-3 \gamma}{4}\right) R_{B}$.
Bondi time: $\quad t_{B}=\frac{R_{B}}{c_{s}}=\frac{G M_{B H}}{c_{s, \infty}^{3}}$.

## Simplest Case: Spherical Bondi Accretion Flow onto a SMBH

- GADGET-3: 3-d cosmological SPH/ N -body code (Springel '05)
- Central SMBH $10^{8}$ M。 represented by a pseudoNewtonian Paczynsky \& Wiita (1980) potential
- $r_{\text {out }}=5-20 \mathrm{pc}, \mathrm{N}_{\text {ptcl }}=64^{3}-128^{3}$
- Set IC to uniform/spherical Bondi flow $w / \gamma=1.01, \rho_{\infty}=10^{-19}$ $\mathrm{g} / \mathrm{cm}^{3}, \mathrm{~T}_{\infty}=10^{7} \mathrm{~K}, \mathrm{~T}_{\text {init }}=\mathrm{T}_{\infty}$
- Corresponding Bondi solution: $\mathrm{R}_{\mathrm{B}}=3 \mathrm{pc}, \mathrm{R}_{\text {sonic }}=1.5 \mathrm{pc}, \mathrm{t}_{\mathrm{B}}=7.9 \mathrm{e} 3 \mathrm{yr}$
- All runs: $\boldsymbol{r}_{\text {in }}=\mathbf{0}$. $1 \mathrm{pc}, \gamma=1.0 \mid$

3-d spherical volume, vacuum
boundary condition

| Run <br> No. | $\begin{gathered} r_{\mathrm{out}} \\ {[\mathrm{pc}]} \end{gathered}$ | $N$ b | IC | $\begin{gathered} M_{\mathrm{tot}, \mathrm{IC}}{ }^{\mathrm{c}} \\ {\left[M_{\odot}\right]} \end{gathered}$ | $\begin{gathered} M_{\text {part }}{ }^{\mathrm{d}} \\ {\left[M_{\odot}\right]} \end{gathered}$ | $\begin{gathered} t_{\mathrm{end}}{ }^{\mathrm{e}} \\ {\left[10^{4} \mathrm{yr}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 5 | $64^{3}$ | Uniform ${ }^{\text {i }}$ | $3.96 \times 10^{5}$ | 1.51 | 3 |
| 2 | 10 | $64^{3}$ | Uniform | $6.19 \times 10^{6}$ | 23.61 | 7.2 |
| 3 | 50 | $128^{3}$ | Uniform | $7.73 \times 10^{8}$ | 368.60 | 20 |
| 4 | 5 | $64^{3}$ | Bondi ${ }^{\text {j }}$ | $1.81 \times 10^{6}$ | 6.89 | 2 |
| 5 | 10 | $64^{3}$ | Bondi | $9.76 \times 10^{6}$ | 37.23 | 8 |
| 6 | 10 | $128^{3}$ | Bondi | $9.76 \times 10^{6}$ | 4.65 | 8 |
| 7 | 20 | $128^{3}$ | Bondi | $6.24 \times 10^{7}$ | 29.75 | 8 |
| $7 a^{\mathrm{k}}$ | 20 | $128^{3}$ | Bondi | $6.24 \times 10^{7}$ | 29.75 | 80 |
| $7 b^{\text {l }}$ | 20 | $128^{3}$ | Bondi | $6.24 \times 10^{7}$ | 29.75 | 100 |
| 8 | 50 | $128^{3}$ | Bondi | $8.48 \times 10^{8}$ | 404.35 | 16 |
| 9 | 20 | $128^{3}$ | $\rho_{B}, v_{\text {init }}=0$ | $6.24 \times 10^{7}$ | 29.75 | 8 |
| 10 | 20 | $128{ }^{3}$ | Uniform | $4.95 \times 10^{7}$ | 23.60 | 8 |
| 11 | 20 | $128^{3}$ | Hernquist ${ }^{\text {m }}$ | $6.24 \times 10^{7}$ | 29.75 | 7.2 |
| $12^{\mathrm{n}}$ | 20 | $128^{3}$ | Bondi | $6.24 \times 10^{7}$ | 29.75 | 8 |

## Example: Properties of Particles

## Run 7:

- $\mathrm{r}_{\text {out }}=20 \mathrm{pc}, \mathrm{N}_{\mathrm{ptcl}}=128^{3}$
- Snap at $\mathrm{t}=2 \mathrm{t}_{\mathrm{B}}=1.6 \mathrm{e} 4 \mathrm{yr}$
- Follows the Bondi solution (red curve) well except the very inner part
- Inner part: supersonic (M~6), outerpart: subsonic








## Mass Inflow Rates at $\mathbf{r}_{\text {in }}$

- the larger rout, the longer duration of Bondi inflow rate
- If started from a Bondi flow, Bondi rate is achieved quickly.
- After a while, the inflow rate decreases due to the artificial outflow at the outer boundary.
- Greater sim. volume reduces this effect on mass inflow.



## Radiative Heating \& Cooling

- Xray emitting corona irradiates the accretion flow

$$
L_{X}=f_{X} L_{\mathrm{Edd}}, \quad L_{\mathrm{Edd}}=\frac{4 \pi c G m_{p} M_{B H}}{\sigma_{e}}, \quad \text { Flux: } \quad F_{X}=\frac{L_{X}}{4 \pi r^{2}} .
$$

- Approx. analytic heating/cooling rates from Blondin '94; opt-thin gas illuminated by a 10 keV bremsstrahlung.

$$
\text { net rate: } \quad \rho \mathcal{L}=n^{2}\left(G_{\text {Compton }}+G_{X}-L_{b, l}\right) \quad\left[\mathrm{erg} \mathrm{~cm}^{-3} \mathrm{~s}^{-1}\right],
$$

Compton h/c rate: $\quad G_{\text {Compton }}=8.9 \times 10^{-36} \xi\left(T_{X}-4 T\right) \quad\left[\mathrm{erg} \mathrm{cm}^{3} \mathrm{~s}^{-1}\right]$.
Net Xray photoioniz. heating and recomb. cooling rate:

$$
G_{X}=1.5 \times 10^{-21} \xi^{1 / 4} T^{-1 / 2}\left(1-\frac{T}{T_{X}}\right) \quad\left[\operatorname{erg~cm}{ }^{3} \mathrm{~s}^{-1}\right]
$$

$$
L_{b, l}=3.3 \times 10^{-27} T^{1 / 2}
$$

$$
+\left[1.7 \times 10^{-18} \exp \left(-1.3 \times 10^{5} / T\right) \xi^{-1} T^{-1 / 2}\right.
$$

$$
\left.+10^{-24}\right] \delta \quad\left[\mathrm{erg} \mathrm{~cm}^{3} \mathrm{~s}^{-1}\right]
$$

$$
\mathrm{T}_{\mathrm{X}}=\mathrm{I} .16 \times 10^{8} \mathrm{~K} \quad(=10 \mathrm{keV}, \text { Blondin '94) }
$$

## Runs with radiative cooling/heating

| Run <br> No. | $\begin{gathered} r_{\mathrm{out}} \\ {[\mathrm{pc}]} \end{gathered}$ | $N$ | $\begin{gathered} M_{\mathrm{tot}, \mathrm{IC}} \\ {\left[M_{\odot}\right]} \end{gathered}$ | $\begin{gathered} M_{\text {part }} \\ {\left[M_{\odot}\right]} \end{gathered}$ | $\gamma_{\text {init }}$ | $\begin{gathered} T_{\infty} \\ {[\mathrm{K}]} \end{gathered}$ | $\begin{aligned} & R_{B} \\ & {[\mathrm{pc}]} \end{aligned}$ | $\begin{gathered} \rho_{\infty} \\ {\left[\mathrm{g} / \mathrm{cm}^{3}\right]} \end{gathered}$ | $T_{\text {init }}$ | $\begin{gathered} L_{X} \\ {\left[L_{\mathrm{Edd}}\right]} \end{gathered}$ | $\begin{gathered} t_{\mathrm{end}} \\ {\left[10^{5} \mathrm{yr}\right]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 13 | 20 | $128^{3}$ | $5.81 \times 10^{5}$ | 0.277 | 1.4 | $10^{7}$ | 2.19 | $10^{-21}$ | $T_{\infty}$ | 0.5 | 1.0 |
| 14 | 50 | $128^{3}$ | $8.23 \times 10^{6}$ | 3.92 | 1.4 | $10^{7}$ | 2.19 | $10^{-21}$ | $T_{\infty}$ | 0.5 | 2.9 |
| 15 | 20 | $128^{3}$ | $5.81 \times 10^{-1}$ | $2.77 \times 10^{-7}$ | 1.4 | $10^{7}$ | 2.19 | $10^{-27}$ | $T_{\infty}$ | 0.5 | 1.0 |
| 16 | 20 | $256^{3}$ | $5.81 \times 10^{-1}$ | $3.46 \times 10^{-8}$ | 1.4 | $10^{7}$ | 2.19 | $10^{-27}$ | $T_{\infty}$ | $5 \times 10^{-4}$ | 1.9 |
| 17 | 20 | $128^{3}$ | $5.81 \times 10^{5}$ | 0.277 | 1.4 | $10^{7}$ | 2.19 | $10^{-21}$ | $T_{\text {rad }}{ }^{\text {b }}$ | $5 \times 10^{-4}$ | 2.9 |
| 18 | 20 | $128^{3}$ | $5.65 \times 10^{5}$ | 0.269 | 5/3 | $10^{7}$ | 1.84 | $10^{-21}$ | $T_{\text {rad }}$ | $5 \times 10^{-4}$ | 3.0 |
| 19 | 20 | $128^{3}$ | $1.47 \times 10^{7}$ | 7.0 | $5 / 3$ | $10^{5}$ | 183.9 | $10^{-21}$ | $T_{\text {rad }}$ | $5 \times 10^{-4}$ | 1.5 |
| 20 | 200 | $256{ }^{3}$ | $1.33 \times 10^{9}$ | 79.09 | 5/3 | $10^{5}$ | 183.9 | $10^{-21}$ | $T_{\text {rad }}$ | $5 \times 10^{-4}$ | 6.5 |
| 21 | 200 | $256^{3}$ | $4.95 \times 10^{8}$ | 29.50 | 5/3 | $10^{7}$ | 1.84 | $10^{-21}$ | $T_{\text {rad }}$ | $5 \times 10^{-4}$ | 8.7 |
| 22 | 200 | $128^{3}$ | $1.33 \times 10^{7}$ | 6.33 | 5/3 | $10^{5}$ | 183.9 | $10^{-23}$ | $T_{\text {rad }}$ | $5 \times 10^{-4}$ | 70 |
| 23 | 200 | $256^{3}$ | $1.33 \times 10^{7}$ | 0.791 | 5/3 | $10^{5}$ | 183.9 | $10^{-23}$ | $T_{\text {rad }}$ | $5 \times 10^{-4}$ | 20 |
| $24^{\text {c }}$ | 200 | $1.24 \times 10^{7}$ | $9.77 \times 10^{6}$ | 0.791 | 5/3 | $10^{5}$ | 183.9 | $10^{-23}$ | $T_{\text {Run23 }}$ | $5 \times 10^{-5}$ | 19 |
| 25 | 200 | $1.24 \times 10^{7}$ | $9.77 \times 10^{6}$ | 0.791 | $5 / 3$ | $10^{5}$ | 183.9 | $10^{-23}$ | $T_{\text {Run23 }}$ | $5 \times 10^{-3}$ | 21 |
| 26 | 200 | $1.24 \times 10^{7}$ | $9.77 \times 10^{6}$ | 0.791 | $5 / 3$ | $10^{5}$ | 183.9 | $10^{-23}$ | TRun23 | $1 \times 10^{-2}$ | 22 |
| 27 | 200 | $1.24 \times 10^{7}$ | $9.77 \times 10^{6}$ | 0.791 | $5 / 3$ | $10^{5}$ | 183.9 | $10^{-23}$ | TRun23 | $2 \times 10^{-2}$ | 25 |
| 28 | 200 | $1.24 \times 10^{7}$ | $9.77 \times 10^{6}$ | 0.791 | 5/3 | $10^{5}$ | 183.9 | $10^{-23}$ | $T_{\text {Run23 }}$ | $5 \times 10^{-2}$ | 50 |

## Ptcl properties w/ radiative heating \& cooling

Representative run:

- red: free-fall scaling
- blue:ZEUS-2d result
- Near the inner radius, excess heating by artificial viscosity is seen.
- Inflow rate is enhanced above Bondi rate, due to lower gas temp: $T$ (rout) $<10^{5} \mathrm{~K}, \mathrm{~T}_{\infty}=10^{5} \mathrm{~K}$


Run \#23: $\mathrm{t}=1 \mathrm{Myr}, \mathrm{Lx}_{\mathrm{x}}=5 \mathrm{e}-4 \mathrm{~L}_{\text {Edd }}, \gamma=5 / 3$ $r_{\text {out }}=200 \mathrm{pc}, 256^{3} \mathrm{ptcls}$




green: free-fall scaling w/ only adiabatic term
$\mathrm{T}_{\mathrm{ff} \text {,ar: }}$ solving internal energy eq. w/ both radiative \& adiabatic term

## Impact of varying Lx on inflow rates

- Restart Run \#23 at t=1.4 Myr, Lx/Ledd=5e-4 orig.
- Runs 24-28: increase Lx
- Dramatic decrease in $\dot{M}_{i n}$ at $L x / L_{\text {Edd }}>0.01$ transition from net inflow to net outflow


Thermal instability due to rad. feedback

## Non-spherical outflow: Run 26: rout $=200 \mathrm{pc}, \mathrm{Lx} / \mathrm{L}_{\mathrm{Edd}}=0.0 \mathrm{I}$

due to rad. feedback


## Ptcl Properties: impact of rad feedback

$$
\begin{gathered}
\text { Run 26: } \text { rout }=200 \mathrm{pc}, \\
L^{2} / \mathrm{L}_{\mathrm{Edd}}=0.0 \mathrm{O}, \mathrm{t}=2.0 \\
\mathrm{Myr}
\end{gathered}
$$

Large scatter due to thermal instability --cold inflow and hot




## outflow




Photoionization parameter:

$$
\xi \equiv \frac{4 \pi F_{X}}{n}=\frac{L_{X}}{r^{2} n},
$$



Outflowing gas near outer BC

## Time Evolution of a Single Ptcl

Run 26: rout $=200 \mathrm{pc}$, $L x / L_{\text {Edd }}=0.0 \mathrm{I}, \mathrm{t}=2.0$ Myr

- Start (triangle): $\mathrm{r}=53 \mathrm{pc}$, $\mathrm{t}=1.4 \mathrm{Myr}$
- End (square): r=lpc, t=l. 8 Myr
-     + symbol: dt=0.004 Myr







## Non-spherical outflow: Run $27: r_{\text {out }}=200 \mathrm{pc}, \mathrm{Lx} / L_{\text {Edd }}=0.02$

due to rad. feedback
Temperature


## Ptcl Properties: impact of rad feedback

Run 27: rout $=200 \mathrm{pc}$, Lx/LEdd $=0.02$






## Non-spherical outflow: <br> Run 28: $r_{\text {out }}=200 p c, L x / L_{\text {Edd }}=0.05$

due to rad. feedback

Temperature
gamma_5by3 / Run37 / LxByLedd_5e-2-IC_snap70 (1.4 - 5 Myr)


Density

$\pm 200 p c$

## Run 28


${ }_{20} \log T$


## Conclusions

- GADGET-3 SPH code can reproduce the spherical Bondi accretion rate properly, but with some limitations.
- spurious heating by Artificial Viscosity near $r_{\text {in }}$ \& artificial outflow at rout due to outer BC are problems for SPH.
- non-spherical in/outflow develops due to rad. feedback via thermal instability, even in the simplest situation that we studied --- connection with NLR? (Paper II)
- Future work: include rad. pressure, rotation, diff geometry, comparison w/ NLR obs, connect with cosmological sim

Run 26: $\mathrm{rout}=200 \mathrm{pc}, \mathrm{Lx} / \mathrm{L}_{\text {edd }}=0.01$

$\pm 30 \mathrm{pc}$ range $\quad(\mathrm{t}=2.047 \mathrm{Myr})$

colder, denser filament-like structures due to non-spherical fragmentation

Run 26: $r_{\text {out }}=200 \mathrm{pc}, L x / L_{\text {Edd }}=0.01$
Zoom-in: inner 4 pc


uo!fDz!uo!otoud 6ol


## Run 27

$r_{\text {out }}=200 p c$,
$L x / L_{E d d}=0.02$

## $\mathrm{t}=1.86 \mathrm{Myr}$

## $\mathrm{t}=2.12 \mathrm{Myr}$

## $\mathrm{t}=2.46 \mathrm{Myr}$

$\log T$

$\log \rho_{\text {gas }}$


